

# **Algebraic Geometry at Sea**

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## **Solitons**



The first description of a soliton was by the Scottish engineer John Scott Russell in 1834: he was conducting experiments on boats on a channel when he observed

" a mass of water [...] assuming the form of a large solitary elevation, a rounded, smooth and well-defined heap of water, which continued its course along the channel apparently without change of form or diminution of speed. "



Figure: A reconstruction of Russell's observation in 1995 (Heriot-Watt University).



Since this first sighting, solitary waves have been source of fascination among scientists, and they are still a very active topic of research.

They were described in mathematical terms by Boussinesq (1877) and Korteveg and de Vries (1895), as particular solutions of the Korteweg - de Vries equation:

$$4u_t - 6uu_x - u_{xxx} = 0$$

This is a famous nonlinear partial differential equation, which describes more generally the motion of waves in a one-dimensional shallow channel, such as the one of Russell.

### A not so shallow channel





Figure: The Bosphorus strait this morning.



Let's look for translation waves solutions to the KdV equation: something of the form

$$u(x,t)=v(x-a\cdot t)$$

where v is a very nice function, for example analytic. Imposing that this is a solution we find

$$-4a\cdot v'-6v\cdot v'-v'''=0.$$

Integrating we get

$$-4a \cdot v - 3(v)^2 - v'' + b = 0.$$

We can multiply by v' to obtain

$$-4a \cdot vv' - 3(v)^2v' - v''v' + b \cdot v' = 0,$$

which we can integrate again to obtain

$$-2a \cdot v^2 - v^3 - \frac{1}{2}(v')^2 + b \cdot v + c = 0$$



To summarize, we need to solve the differential equation

$$(v')^2 = -4v^3 - 4a \cdot v^2 + 2b \cdot v + 2c$$

Up to a constant, this is the same as the famous Weierstrass equation:

$$y^2 = f(x),$$
  $f(x) = -4x^3 - 4a \cdot x^2 + 2b \cdot x + 2c$ 

Geometrically this cuts out an elliptic curve  $E = \{(x, y) | y^2 = f(x)\}$  in  $\mathbb{C}^2$  and we are looking for an analytic function  $v : \mathbb{C} \to \mathbb{C}$  such that we have a well-defined map

$$\mathbb{C} \longrightarrow E \qquad z \mapsto (v(z), v'(z))$$

But such a function is known since more than 100 years: it is the Weierstrass  $\wp$ -function, or, equivalently, the (second logarithmic derivative of the) theta function of the curve *E*.

This discussion works much more generally.

## **Algebraic curves and Riemann surfaces**

An algebraic curve is a variety of dimension one, the zero locus of polynomial equations in  $\mathbb{C}^n$  or  $\mathbb{P}^n$ .

$$\left\{x^4+y^4=1\right\}$$



Smooth curves correspond to Riemann surfaces: complex manifolds of dimension one.







Algebraic curves are connected to the Kadomtsev-Petviashvili (KP) equation:

$$\frac{\partial}{\partial x}\left(4u_t-6uu_x-u_{xxx}\right)=3u_{yy}$$

which describes the evolution of two-dimensional waves in shallow water.

This is a generalization of the KdV equation

$$4u_t - 6uu_x - u_{xxx} = 0$$



Figure: KP equation in French waters (Wikipedia).



The connection goes through Riemann's theta function. Integrating holomorphic differentials on a genus g curve, we obtain the Riemann matrix:

$$au$$
 " = "  $\left(\int_{b_j}\omega_i\right)$ 

which is a complex symmetric  $g \times g$  matrix with positive definite imaginary part.

**Riemann's theta function** The associated theta function is a function of  $z \in \mathbb{C}^{g}$  given by

$$\theta(z,\tau) = \sum_{n\in\mathbb{Z}^g} \mathbf{e}\left(\frac{1}{2}n^t\tau n\right) \cdot \mathbf{e}(n^tz) = \sum_{n\in\mathbb{Z}^g} a_n(\tau) \cdot \mathbf{e}(n^tz).$$

where  $e(x) = exp(2\pi i x)$ . This is an infinite sum of exponentials, and it is computable: Maple, MATLAB, Sage, Julia,...



In geometric terms, the Riemann matrix  $\tau$  corresponds to an abelian variety:

$$\mathsf{A}_{ au} = \mathbb{C}^{\mathsf{g}} / (\mathbb{Z}^{\mathsf{g}} + au \mathbb{Z}^{\mathsf{g}})$$

We can construct an abelian variety starting from any Riemann matrix  $\tau$ . If this comes from a curve, the abelian variety is called the Jacobian of the curve *C*.





### **Curves and waves**



Krichever proved that the KP equation has a solution of the form

$$u(x, y, t) = 2 rac{\partial^2 \log}{\partial^2 x} heta(\mathbf{U}x + \mathbf{V}y + \mathbf{W}t, au)$$

for certain vectors  $\mathbf{U}, \mathbf{V}, \mathbf{W} \in \mathbb{C}^{g}$ .

Shiota then solved the Schottky problem: a matrix  $\tau$  represents a Jacobian if and only if the above is a solution of the KP equation.

The first result is very concrete: on the right, a wave obtained from a quartic plane curve in Julia.





- Classify the KP solution arising from a fixed curve: the Dubrovin Threefold. Together with T. Çelik and B. Sturmfels.
- What happens when the curve degenerates to a singular nodal curve? Tropical geometry. Together with C. Fevola, Y. Mandelshtam and B. Sturmfels.
- What happens when the curve degenerates to a cuspidal curve? Rational solutions. Together with T. Çelik and J. Little.
- Try to recover the curve from the matrix τ: Effective Torelli theorem. Together with T. Çelik and D. Eken.



One proof of Krichever's theorem is based on two essential ingredients:

- 1. The algebraic structure of the KP equation: **Sato Grassmannian** of Sato, Segal-Wilson.
- 2. Abel's theorem and Riemann-Roch.

In particular, it can be replicated on singular curves as well: since Abel's theorem and Riemann- Roch still hold.

### **Example: soliton solutions**



Consider the nodal cubic  $C = \{(x, y) | y^2 = x^3 + x^2\} \subseteq \mathbb{C}^2$ . This is the image of

$$\mathbb{C} \longrightarrow C$$
,  $t \mapsto \left(t^2 - 1, t^3 - t\right)$ 

A basis of canonical differentials is given by  $\omega = \left(\frac{1}{t-1} - \frac{1}{t+1}\right) dt$ , and  $\int_{\gamma} \omega = 2\pi i$ . Hence

$$\mathsf{Jac}(C) = \mathbb{C}/2\pi i\mathbb{Z} \stackrel{\sim}{\longrightarrow} \mathbb{C}^*, \quad z \mapsto \exp(z)$$

The theta function in this case is

$$heta(z) = \exp(z) - 1$$

This is a **finite** sum of exponentials, and gives rise to a **soliton solution**.

Note that the nodal cubic  $y^2 = x^3 + x^2$  is a special, or degenerate, case of the cubic  $y^2 = f(x)$ .



Fevola, Mandelstham, Sturmfels and myself, studied these degenerations and the corresponding solitons, from the point of view of tropical geometry.

**Theorem (A., Fevola, Mandelstham, Sturmfels - 2021; Ichikawa - 2023)** Under the degeneration to a rational nodal curve, the theta function becomes a finite sum of exponentials:

$$\theta(z) = \sum_{n \in \mathbb{Z}^g} \exp\left(\pi i n^T \tau n + 2\pi i n^T z\right) \rightsquigarrow \hat{\theta} = \sum_{c \in \mathcal{C}} a_c \exp(c^T z)$$

The corresponding solution of the KP equation is a soliton solution.

### **Example: rational solutions**



Consider the cuspidal cubic  $C = \{(x, y) | y^2 = x^3\} \subseteq \mathbb{C}^2$ . This is the image of

$$\mathbb{C} \longrightarrow C$$
,  $t \mapsto (t^2, t^3)$ 

A basis of canonical differentials is given by  $\omega = \frac{1}{t^2} dt$ , and  $\int_{\gamma} \omega = 0$ . Hence

$$\mathsf{Jac}(\mathcal{C})=\mathbb{C}$$

The theta function in this case is

$$\theta(z)=z.$$

This is a polynomial, and gives rise to a **rational solution**.

Note that the cuspidal cubic  $y^2 = x^3$  is an even more degenerate, case of the cubic  $y^2 = f(x)$ .

# **Polynomial theta functions**



Together with Çelik and Little, we have classified which curves give rise to polynomial theta functions.

**Theorem (A., Çelik, Little - 2021)** The theta function of *C* is polynomial if and only if the curve is rational with unibranched singularities. Moreover, the resulting polynomial is of degree at most  $\frac{g(g+1)}{2}$  and gives a rational solution to the KP equation.

In general we have:

Curve	Theta function	<b>KP</b> solution
smooth	infinite sum of exponentials	quasiperiodic
nodal rational	finite sum of exponentials	soliton
unibranched rational	polynomial	rational



- When do we get real soliton solutions out of rational nodal curves? Tropical geometry and positive geometry. Work in progress by S. Abenda, T. Çelik, C. Fevola, Y. Mandelshtam.
- When do we get real rational solutions out of rational unibranched curves? Discussion in progress with T. Çelik and J. Little.
- What about real real solutions?



#### What if we want to actually model water waves this way? A natural idea:



J. Hammack, N. Scheffner and H. Segur



(a) (d) (c)

# Figure: from Hammack, Scheffner, Segur, Journal of Fluid Mechanics (1989)

Figure: from Dubrovin, Flickinger, Segur, Studies in applied mathematics (1997)

# And today?



In a joint work in progress with M. Bennett, T. Çelik, B. Deconinck and C. Wang, we consider the problem of reconstructing the data of a Riemann matrix  $\tau$ , together with the wave vectors **U**, **V**, **W**, from one KP solution

$$u(x, y, t) = 2 \frac{\partial^2 \log}{\partial^2 x} \theta(\mathbf{U}x + \mathbf{V}y + \mathbf{W}t, \tau)$$

This was already done by Dubrovin, Flickinger and Segur, but only with genus two solutions, and with techniques valid exclusively in genus two.

# An algorithm



Suppose *f* is a periodic function of *x* with period *U*, then we have

$$f(x) = \sum_{k \in \mathbb{Z}} c_k \cdot \exp(2\pi i \cdot k \cdot Ux), \quad \hat{f}(\xi) = \sum_{k \in \mathbb{Z}} c_k \delta(\xi - 2\pi Uk).$$

So, if the coefficients  $c_k$  decay sufficiently fast, the Fourier transform of f has "peaks" around  $2\pi k$ , for small values of k.

- ▶ We look the Fourier transform of u(x, 0, 0). This will have some "peaks", which turn out to be the coordinates of **U**. Same with y, z and **V**, **W** respectively.
- We can recover the function  $\tau(x, y, z) = \theta(\mathbf{U}x + \mathbf{V}y + \mathbf{W}z, B)$  from U.
- Recovering the matrix  $\tau$  becomes a matter of numerical linear algebra:

$$au(x,0,0) = \sum_{n\in\mathbb{Z}^g} \mathbf{e}\left(\frac{1}{2}n^t \tau n\right) \cdot \mathbf{e}(n^t \mathbf{U}x)$$



# Thank you!